Statistical Mechanics of Two-Dimensional Shuffled Foams: Prediction of the Correlation between Geometry and Topology

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We propose an analytical model for the statistical mechanics of shuffled two-dimensional foams with moderate bubble size polydispersity. It predicts without any adjustable parameters the correlations between the number of sides \( n \) of the bubbles (topology) and their areas \( A \) (geometry) observed in experiments and numerical simulations of shuffled foams. Detailed statistics show that in shuffled cellular patterns \( n \) correlates better with \( \sqrt{A} \) (as claimed by Desch and Feltham) than with \( A \) (as claimed by Lewis and widely assumed in the literature). At the level of the whole foam, standard deviations \( \Delta n \) and \( \Delta A \) are in proportion. Possible applications include correlations of the detailed distributions of \( n \) and \( A \), three-dimensional foams, and biological tissues.

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Cellular materials are interesting both as disordered media with well-defined structural elements, and as models for more complex systems such as biological tissues [1]. Among them, foams are ubiquitous in our daily lives and in many industries [2–4]. Bubble monolayers are much easier to observe and to study (Fig. 1). Such quasi-two-dimensional foams are characterized by their number of bubbles \( N \), area distribution \( p(A) \), and number-of-sides distribution \( p(n) \). Their bubbles are polygonal, with shapes that are locally governed by Plateau’s laws [3,4]. Each side is a thin liquid film with a uniform surface tension; its curvature is determined by the difference of pressure between the two bubbles it separates.

Bubble size distribution and packing (or “topology”) are crucial in determining, e.g., rheological properties [4]. Statistically, a relatively large bubble, within a given foam sample, has more neighbors than a smaller one. To quantify this intuitive impression, bubbles which have the same number \( n \) of neighbors can be grouped, and their normalized average size plotted versus \( n \). One debate [5–14] concerns whether \( n \) correlates with \( A \) or \( \sqrt{A} \) [Figs. 2(a) and 2(b) and Eqs. (1) and (5)]. A second debate [10,11,13] asks the following: can we understand these empirical correlations, and prove them; what is their physical origin, and why do they apply only to some cellular patterns? We ask the same questions about empirical observation [15] of proportionality between the geometrical disorder, or area polydispersity, \( \Delta A \) = \( \sqrt{A^2} - \langle A^2 \rangle / \langle A \rangle \), and the topological disorder, \( \Delta n \) = \( \sqrt{n^2} - \langle n^2 \rangle / \langle n \rangle \), where \( \langle \cdot \rangle \) denotes the average over all bubbles in a foam and \( \Delta \) the standard deviation [Fig. 2(c)].

Existing models [2,8,10,11,16–19] seek distributions of \( n \) and \( A \) which are optimal in some sense. This is relevant at times long enough for \( p(A) \) to vary, e.g., under coarsening or bubble coalescence [2–4]. At shorter time scales, \( N \) and \( p(A) \) are fixed: only the number of sides of the bubbles varies, when they undergo “T1” neighbor changes [Fig. 1(d)]. At most instants the foam is in a local energy minimum, that is, a metastable state; each T1 is instantaneous and followed by a relaxation towards a new local energy minimum [4]. Local minima have almost the same energy, and are separated by high energy barriers corresponding to T1s [17]. There is no process leading towards a global energy minimum. T1s instead induce a random exploration of local minima, effectively “shuffling” the foam [15]. This notion of shuffling is empirically defined [15,16] as having had enough T1s per bubble to forget about the foam’s initial preparation. Numerical simulations and theoretical analyses have shown that cycles of shear of amplitude significantly larger than the yield strain almost completely remove (although seldom perfectly) both residual trapped stresses and spurious correlations [20].

In the present Letter, we obtain, for shuffled foams of moderate polydispersity, an analytical expression for the distribution of \( n \) for any given distribution of \( A \). We use a grand-canonical description [21] with a constraint based on mechanical equilibrium and space filling (the curvature sum rule [2–4]) rather than on energy. Predictions of geometry-topology correlations agree with shuffled foam data obtained from existing experiments on liquid foams [15] and new, refined, numerical simulations of two sorts.
Surfaces evolve (SE) simulations represent each side as an arc of a circle [Fig. 1(c)]; they provide detailed information about each bubble’s position and shape. Each foam is prepared from a Voronoi construction based on seed points generated by a random Poisson point process [26,27], with the bubble areas adjusted to fit the desired area distribution where necessary, before convergence to an energy minimum. We cyclically shear, by deforming the shape of the periodic box, about 400 foams with \( N = 2500 \) bubbles in two perpendicular directions with strain amplitude 1.5.

In experiments and simulations we made bidisperse foams, with the ratio of large to small bubble areas ranging from 1 to 8, and polydisperse foams with normal (Potts), Poisson (SE), or intermediate (experiments) distributions of areas, with relative width \( \Delta A/A \) ranging from 0 to 1.15.

The topology of the foam clearly depends on the relative (and not absolute) bubble sizes. To enable comparisons with simulations and theory, in what follows all data are plotted as dimensionless. We first present averages \( \bar{n} \) over individual bubbles. Figure 2(a) shows the intuitively expected increase of \( \bar{n} \) with the bubble size. Moreover, it shows that all simulation data of \( \bar{n} \) vs \( \sqrt{A}/(\sqrt{A}) \) are grouped, and display a linear variation, on a much wider range than \( \bar{n} \) vs \( A/(\sqrt{A}) \). These ensemble averages require that we bin together bubbles of similar sizes. In addition, in SE simulations we track for each bubble the fluctuations of \( A \) with time (\( A \) being fixed) and check (not shown here) that time averages yield exactly the same dependence of \( \bar{n} \) on \( A \) as ensemble averages.

Second, representations of bubble sizes averaged over bubbles of the same \( n \), without binning, are easy to extract from experiments or simulations, which probably explains their prevalence in the literature [5–11]. Figure 2(b) shows the intuitively expected increase of \( \sqrt{A}/(\sqrt{A}) \) with \( n \), but here without obvious superposition of the data; plots (not shown) with bubble perimeter \( P \) rather than \( \sqrt{A} \) are similar. The fit is linear; its slope is very different from the inverse of the slope in Fig. 2(a), due to the statistical character of the relationship between \( \sqrt{A} \) (or \( A \)) and \( n \) [13].

Third, Fig. 2(c) shows measures of the foam’s disorder. The topological disorder is proportional to the geometrical one up to \( \sim 0.4 \). At higher polydispersity, measures of disorder are less correlated.

Our model for these data is based on the ideas introduced in [21], which we now briefly recall. Each T1 keeps two quantities constant: the sum of \( n \) over the four bubbles involved, and the sum of the four bubbles’ curvatures \( \kappa \) [28]. A foam’s macrostate is thus defined by its total curvature (summed over \( N \) bubbles) \( \kappa_{\rm tot} = \sum \kappa \) and total number of sides \( N_s = \sum n \). Within a macrostate, the set of microstates is defined as the set of all accessible local energy minima: we assume they all have the same probability, and use a grand-canonical description:

\[ \Delta n/n \rightarrow 0 \]
consider one particular bubble of area $A$, the rest of the foam constitutes a reservoir of sides and curvature, exchanged through $T_1$'s. We then use a mean field approximation, disregarding neighboring correlations and assuming this reservoir foam is homogeneous and isotropic. The probability for a bubble with size $A$ to have $n$ sides is

$$p_A(n) = \frac{A}{\sqrt{\pi}} e^{-(n-A)^2/(2\sigma^2)}$$

where $A = \sum_{n=1}^{\infty} e^{-(n-A)^2/(2\sigma^2)}$ is the partition function of the entire foam, defined through $\ln \Xi = N \int_0^\infty p(A) \ln \chi(A) dA$. For a very large foam ($N \to \infty$), the constraint of space filling sets the values ($\kappa_{tot}$) through $\langle \kappa_{tot} \rangle = \sum_{n \geq 1} e^{-(n-A)^2/(2\sigma^2)}$ and $\langle N_c \rangle = \sum_{n \geq 1} e^{-(n-A)^2/(2\sigma^2)}$, where $\Xi$ is the partition function of the entire foam. To proceed, we obtain the values of $\beta$ and $\mu$ by solving the system:

$$\frac{\partial \ln \Xi}{\partial \beta} = 6N$$

and

$$\frac{\partial \ln \Xi}{\partial \mu} = 0.$$
assumption \( \sigma \ll \bar{n}(A) - 3 \) requires \( A^{1/2}/(\langle A^{1/2} \rangle) \gg \langle A^{1/2}\rangle(A^{-1/2} - 1) \). Thus, the criterion on the width of \( p(A) \) is \( \langle A^{1/2}\rangle(A^{-1/2} - 1) \ll 1 \). It is more intuitive to characterize the width of \( p(A) \) by its normalized standard deviation \( \Delta A/\langle A \rangle \). For that purpose, we change variable and use the relative deviation \( e \) from monodispersity [30], assumed to be \( \ll 1 \): the condition of validity of this approximation, and thus of Eq. (1), reduces to \((\Delta A/\langle A \rangle)^2 \ll 4\), and is thus obeyed by both small and moderate dispersions.

Using the same change of variable, \( \bar{n}(A) \) can be expressed as a function of \( A \) [30], instead of \( \sqrt{A} \). For foams with a small area disparity \((\Delta A/\langle A \rangle \ll 1)\) we can write

\[
\bar{n}(A) \approx \frac{3}{2} \left( 3 + \frac{A}{\langle A \rangle} \right) \quad (5)
\]

The inset of Fig. 2(a) shows the linear variation predicted by Eq. (5). Equations (1) and (5) are consistent. As expected, Eq. (1) has a range of validity larger than that of Eq. (5): roughly \( A/\langle A \rangle \in [0.16, 2.56] \) vs \([0.50, 1.50] \). Note that we have no simple prediction for the Desch-Feltham or Lewis representation [Fig. 2(b)].

We can also calculate \( \Delta n/\langle n \rangle \), using \( \langle n^2 \rangle = \int_0^\infty p(A) \chi(A)^{-1}(\sigma^2/\chi) d\mu \) and Eq. (2). This yields \( (\Delta n/\langle n \rangle)^2 = [(\langle A^{1/2} \rangle(A^{-1/2} - 1) + \langle A \rangle(A^{1/2} - 2)/4 \].

For shuffled foams with a moderate disparity we obtain [30]:

\[
\frac{\Delta n}{\langle n \rangle} = \frac{1}{2^{3/2}} \frac{\Delta A}{\langle A \rangle} = 0.35 \frac{\Delta A}{\langle A \rangle}. \quad (6)
\]

Despite the approximations involved in the analytical derivation of Eq. (6), its prediction is indiscernible from the linear fit to the data over the whole expected range, that is, up to a geometrical disorder of 0.4 [Fig. 2(c)]. Note that the agreement is better with experiments than numerics, especially at low disorder. We attribute this small discrepancy to the lower number of bubbles in Potts simulations (\( \sim 2500 \) bubbles for experiments and SE versus 112 for Potts), and to the much lower value of the liquid fraction for SE simulations (\( \sim 10^{-3} - 10^{-2} \) for Potts and experiments versus \( \sim 10^{-5} \) for SE [20]). At larger dispersions the relation between \( \Delta n/\langle n \rangle \) and \( \Delta A/\langle A \rangle \) depends on the exact shapes of the area distributions: this explains the scatter in the data.

In conclusion, in shuffled two-dimensional foams we predict without any free parameter the correlation between topology and geometry: \( n \) correlates better with \( \sqrt{A} \) than with \( A \) [Eq. (1)], and the two measures of disorder are proportional [Eq. (6)]. Although the exact shape of the area distribution plays a role in principle, we expect that in a linear approximation its mean and width play the dominant role. This is what our linear derivation captures. These results should lead to a more accurate description of other shuffled cellular patterns such as some biological tissues and metallic grains, and of shuffling and ergodicity in soft matter.

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References:


[28] The curvature $\kappa$ of a bubble is the sum of the algebraic curvatures of its sides, which we count here as positive if the center of curvature is outside the bubble.


[30] We use the change of variable $A = \langle A \rangle (1 + \epsilon)$. Assuming $\epsilon \ll 1$, we can write any $k$th moment of $A$ as an expansion in $\epsilon$, up to order $O(\epsilon^3)$: $\langle A^k \rangle = \langle A \rangle^k (1 + k(k - 1)\epsilon^2)/2$. Thus we obtain $\langle A^{1/2} \rangle \times \langle A^{-1/2} \rangle - 1 \approx \langle \epsilon^2 \rangle/4$; similarly, $\langle \Delta A / \langle A \rangle \rangle^2 \approx \langle \epsilon^2 \rangle$; and finally $[\langle A^{1/2} \rangle \langle A^{-1/2} \rangle + \langle A \rangle \langle A^{1/2} \rangle^{-2} - 2]/4 \approx \langle \epsilon^2 \rangle/8$. 